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Deriving The Upper Blow-up Rate Estimate for a Parabolic Problem

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Abstract

In this paper, the blow-up solutions for a parabolic problem, defined in a bounded domain, are studied. Namely, we consider the upper blow-up rate estimate for heat equation with a nonlinear Neumann boundary condition defined on a ball in R_n .

Keywords: Blow-up solutions; Parabolic problem; Heat equation ; Neumann boundary condition; outward normal vector.

اشتقاق صيغة لتقدير النسبة العليا للانفجار لمسألة قطع مكافئ

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الخلاصه

في هذا البحث ، تم دراسه الحلول المنفجرة لمسأله قطع مكافئ معرفة على فضاء مقيد. وبصورة ادق، نهتم بالنسبه العليا لتقدير الانفجار لمعادلة الحرارة مع شرط نيومان الحدودي غير الخطي معرفة على كرة في الفضاء Rn .

1. Introduction

In this paper, we study the following parabolic problem:

$$\begin{array}{ll} u_t = \Delta u, & (x,t) \in B_R \times (0,T), \\ \frac{\partial u}{\partial \eta} = \lambda u^p e^{qu} &, & (x,t) \in \partial B_R \times (0,T), \\ u(x,0) = u_0(x), & x \in B_R, \end{array}$$
(1)

where $p \ge 1$; $q, \lambda > 0$; B_R is a ball in R^n ; η is the outward normal and u_0 is smooth, nonzero, nonnegative, radially symmetric, satisfying the following condition:

$$\frac{\partial u_0}{\partial \eta} = \lambda u_0^p e^{q u_0} \quad , \quad x \in \partial B_R,$$
$$\Delta u_0 \ge 0, \quad x \in \overline{B}_R.$$

Many real problems in the fields of fluid dynamics, population, heat propagation, and others are modeled using the forms of parabolic partial differential equations assosiated with different types of initial-boudary conditions [1-3].

In some cases, the solutions of parabolic problems cannot be continued globally in time. This is called the blow-up phenomenon which, in time-dependent problems, has been studied over the past years by many authoes [3-7]. One of these problems is the problem of the heat equation defined in a ball B_R with a nonlinear Neumann boundary condition, $\frac{\partial u}{\partial \eta} = f(u)$ on $\partial B_R \times (0,T)$, which has been introduced in previous articles [7-11].

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In an earlier study [10], it was proved that if 1/f is integrable at infinity for positive values, and f is nondecreasing, then with any positive initial function u_0 , the blow-up occurs in a finite time. In addition, if $f \in C^2(0, \infty)$, is a convex and increasing function in $(0, \infty)$, then the blow-up can only occur on the boundary.

In another study [8], a speial case was considerd, where $f(u) = u^p$, showing that for any u_0 , the blow-up occurs for p > 1, and it can only occur on the boundary. Furthermore, in other investigations [9,11], it was proved that the upper (lower) blow-up rate estimates are as follows:

$$C_1(T-t)^{\frac{-1}{2(p-1)}} \le \max_{x \in \overline{B}_R} u(x,t) \le C_2(T-t)^{\frac{-1}{2(p-1)}}, \quad t \in (0,T).$$

Another special case, where $f(u) = e^u$, was considered in another work [7], showing that all positive solutions blow up in finite times, and the blow-up can only occur on the boundary. Moreover, the upper (lower) blow-up rate estimates are as follows:

$$C_1(T-t)^{-1/2} \le e^{u(1,t)} \le C_2(T-t)^{-1/2}, \quad 0 < t < T.$$

A similer result has been obtained [3] for another case: $f(u) = e^{u^p}$, p > 1. In this paper, we derive that upper blow-up rate estimate for problem (1), showing that:

$$\max_{\overline{B}_R} u(x,t) \le \log C - \frac{\lambda}{2q} \log(T-t), \quad 0 < t < T$$

The rest of this paper is organized as follows: In section two, we discuss the local existance, blowup, and blow-up set with stating the main properties of classical solutions for problem (1). In section three, we derive the upper blow-up rate estimate. Section four is devoted to state some conclusions.

2. Basic properties and Blow-up Set

It is well known that with any smooth initial function, problem (1) has a local unique classical solution [12]. Moreover, since $f(u) = \lambda u^p e^{qu}$ is positive, increasing, convex, and 1/f is integrable at infinity for u > 0, then each solution of problem (1) blows up in a finite time and the blow-up can only occur on the boundary [10].

The next lemma, proved in an earlier study [3], presents some solution properties of problem (1). For simplicity, we denote u(r,t) = u(x,t).

Lemma 2.1. Let $u \in C^{4,2}(B_R \times (0,T))$ be a solution to problem (1). Then

- 1. *u* is positive and radial on $\overline{B}_R \times (0, T)$, and u_r is nonnegative in $[0, R] \times [0, T)$.
- 2. u_t is positive in $\overline{B}_R \times (0,T)$, and if $\Delta u_0 \ge a > 0$, in \overline{B}_R , then $u_t \ge a$, in $\overline{B}_R \times [0,T)$.

3. Upper Blow-up Rate Estimate

The next theorem is concerned with deraving the upper blow-up rate estimate for problem (1). **Theorem 3.2** Let $u \in C^{4,2}(B_R \times (0,T))$ be a blow-up solution to (1), such that $\Delta u_0 \ge a > 0$ $\forall x \in \overline{B}_R$. Then there is a positive constant *C* such that

$$\max_{\overline{B}_R} u(x,t) \le \log C - \frac{1}{2q} \log(T-t), \quad 0 < t < T.$$
⁽²⁾

Proof.

As in another work [5], we define the function:

 $F(x,t) = u_t(r,t) - \varepsilon u_r^2(r,t), \quad (x,t) \in B_R \times (0,T).$ By some direct calculations, it follows that

$$F_t - \Delta F = 2\varepsilon \left(\frac{n-1}{r^2}u_r^2 + u_{rr}^2\right) \ge 0.$$

Since $\Delta u_0 \ge a > 0$, and $u_{0r} \in C(\overline{B}_R)$, for a samll enough value ε , we have $F(x, 0) = \Delta u_0(r) - \varepsilon u_{0r}^2(r) \ge 0$, $x \in B_R$,

In addition,

$$\frac{\partial F}{\partial \eta}\Big|_{x \in S_R} = u_{rt}(R,t) - 2\varepsilon u_r(R,t) u_{rr}(R,t)$$
$$= (\lambda u^p e^{qu})_t - 2\varepsilon \lambda u^p e^{qu} (u_t(R,t) - \frac{n-1}{r} u_r(R,t))$$
$$\geq (qu + n - 2\varepsilon) \lambda u^{p-1} e^{qu} u(R,t)$$

 $\geq (qu + p - 2\varepsilon)\lambda u^{p-1}e^{qu} \quad u_t(R, t).$ Since $u_t > 0$, on $\overline{B}_R \times (0, T)$, it follows that

$$\frac{\partial F}{\partial \eta}|_{x \in S_R} \ge 0, \quad t \in (0,T)$$

provided that

 $\varepsilon \leq \frac{qu_0+p}{2}.$ From the comparison principles [13] and [14], we obtain $F(x,t) \ge 0, \quad (x,t) \in \overline{B}_R \times (0,T),$ In particular, $F(x,t) \ge 0$, for |x| = R, which leads to $u_t(R,t) \ge \varepsilon u_r^2(R,t) = \varepsilon \lambda^2 u^{2p} e^{2qu}$, $t \in (0,T)$. Since u blows up at T, and increases in time, there exists $\tau \leq T$ such that $u(R,t) \ge 1$ for $\tau \le t < T$, which leads to $u_t(R,t) \ge \varepsilon \lambda^2 e^{2qu(R,t)}, \quad t \in [\tau,T).$ By integrating the last inequality from t to T, we obtain $\int_t^T u_t e^{-2qu(R,t)} \ge \varepsilon \lambda^2 (T-t).$ Thus $-\frac{1}{2q}e^{-2qu(R,t)}|_t^T \ge \varepsilon\lambda^2(T-t).$ (3)

Since

$$u(R,t) \to \infty$$
, $e^{-qu(R,t)} \to 0$ as $t \to T_{t}$

So that, from the inequality (3), we get

$$\frac{1}{e^{qu(R,t)}} \ge (2\lambda^2 q\varepsilon(T-t))^{1/2},$$

Thus

$$(T-t)^{\frac{1}{2}}e^{qu(R,t)} \le \frac{\lambda}{\sqrt{2q\varepsilon}}$$

So, for some postiive constant C, we have

$$\max_{\overline{B}_R} u(x,t) \le \log C - \frac{\lambda}{2q} \log(T-t), \quad 0 < t < T.$$

4. Conclusions

This paper is devoted to derive the upper blow-up rate estimate for problem (1). The results show that the upper blow-up rate estimate formula does not depend on p, which means that the power function, appeared in the boundary condition, does not make any effect on the blow-up profile to problem (1). Therefore, the influence of the power function may only appear on the blow-up time.

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